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## On the unboundedness of martingale transforms

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The starting point for our investigation is an observation of Stein and Weiss (1959) or more precisely Davis' (1973) proof of this fact. To state their results and explain our motivation, we will need a number of definitions:

Let $B_{t}$ be a two dimensional Brownian motion.
Let $D=\{z:|z|<1\}$
Let $\tau=\inf \left\{t: B_{t} \notin D\right\}$
Let $E \subset \partial D$ and let $u(x)=P_{x}\left(B_{\tau} \in E\right)$.
Finally let $v(x)$ be the "harmonic conjugate" of $u$ : i.e. the unique function with $v(0)=0$ which makes $u+i v$ an analytic function.

The function $u$ is an object which has been much studied by probabilists (see e.g. Port and Stone (1978), F. Knight (1981), or Chung (1982)) and it is well known that $u$ is harmonic in $D$ and

$$
\begin{equation*}
\lim _{t \uparrow \tau} u\left(B_{t}\right)=1_{E}\left(B_{\tau}\right) \quad \text { a.s. } \tag{1}
\end{equation*}
$$

Stein and Weiss' result shows that u's harmonic conjugate is also special.
(2) $\quad \lim _{t \uparrow \tau} v\left(B_{t}\right)$ exists a.s. and furthermore the distribution of the limit $t \uparrow \tau$ depends only on $P_{0}\left(B_{\tau} \in E\right)$.

Stein and Weiss proved (2) by supposing $E$ was a finite union of intervals and then patiently finding the places where $v\left(e^{i \theta}\right)>y$. See pp. 273-274. In (1973) Davis gave the following proof of their result which makes the conclusion obvious.

Proof of (2). Itô's formula implies that if $t<\tau$

$$
\begin{aligned}
& u\left(B_{t}\right)=\int_{0}^{t} \nabla u\left(B_{s}\right) \cdot d B_{s} \\
& v\left(B_{t}\right)=\int_{0}^{t} \nabla v\left(B_{s}\right) \cdot d B_{s}
\end{aligned}
$$

and the Cauchy Riemann equations:

$$
\frac{\partial v}{\partial x}=-\frac{\partial u}{\partial y} \quad \frac{\partial v}{\partial y}=\frac{\partial u}{\partial x}
$$

imply $\nabla \mathrm{u} \cdot \nabla \mathrm{v}=0$ and $|\nabla \mathrm{u}|=|\nabla \mathrm{v}|$ so it follows from Lévy's theorem (see Meyer
(1976), or Durrett (1984), Section 2.11) that $\left(u\left(B_{t}\right), v\left(B_{t}\right)\right) t<\tau$ is a time change of a Brownian motion $\bar{B}_{u}$ run for a random amount of time $u<\sigma$.

To prove (2) we will show that $\sigma=T \equiv \inf \left\{u: \bar{B}_{u}^{1} \notin(0,1)\right\}$. If we discard the trivial cases $P_{0}\left(B_{\tau} \in E\right)=0$ or 1 then $0<u(x)<1$ for $x \in D$ and hence $u\left(B_{t}\right) \in(0,1)$ for $t<\tau$ so $\sigma \geq T$. (1) shows we cannot have $\sigma>T$ so we must have $\sigma=T$.

To motivate our generalization we begin by redescribing the relationship between $u$ and v. It is well known (see Meyer (1976) or Durrett (1984), Section 2.14) that
(3) If $X \in \sigma\left(B_{t} t \geq 0\right)$ has $E X=0$ and $E X^{2}<\infty$ then

$$
x=\int_{0}^{\infty} H_{s} \cdot d B_{s}
$$

where

$$
E X^{2}=E \int_{0}^{\infty}\left|H_{s}\right|^{2} d s
$$

Let $A$ be a $d \times d$ matrix. Since

$$
E \int_{0}^{\infty}\left|A H_{s}\right|^{2} d s \leq E \int_{0}^{\infty} C\left|H_{s}\right|^{2} d s=C E X^{2}<\infty
$$

the Burkholder Gundy inequalities (see Meyer (1976) or Durrett (1984), Section 6.3) imply that $\int_{0}^{t} A H_{S} \cdot d B_{S}$ is an $L^{2}$ bounded martingale so we can define a new random variable by setting

$$
A * X=\int_{0}^{\infty} A H_{S} \cdot d B_{S}
$$

(see Durrett (1984), Section 6.6 for more details).
$A * X$ is called a martingale transform. If $d=2$ and we let $A=\left(\begin{array}{rr}0 & -1 \\ 1 & 0\end{array}\right)$, $X=1_{(B \in E)}$ then using the notation introduced in the proof of (2) the Cauchy Riemann ${ }^{\tau}$ equations can be written as $\nabla v=A \nabla u$, and it follows that $A * X=v\left(B_{\tau}\right)$.

With conjugation identified as a martingale transform, it becomes natural to ask when (2) holds for martingale transforms. Tracing back through the proof of (2) gives the following result:
(4) Suppose $A$ satisfies (a) $y \cdot A y=0$ and (b) $|y|=|A y|$ for all $y \in R^{d}$ then the distribution of $A * 1_{B}$ depends only on $P(B)$.
Unfortunately matrices which satisfy both (a) and (b) are rare. There are none if $d$ is odd because such matrices must have a real eigenvector and yet
(a) $\Rightarrow$ there is no nonzero real eigenvalue
(b) $\Rightarrow A$ is invertible $\Rightarrow 0$ is not an eigenvalue.

In even dimensions the situation is somewhat better but not much. It is easy to see that there are examples

$$
\left(\begin{array}{rr}
0 & 1  \tag{5}\\
-1 & 0
\end{array}\right),\left(\begin{array}{rr|rr}
0 & 1 & 0 & 0 \\
-1 & 0 & 0 & 0 \\
\hline 0 & 0 & 0 & 1 \\
0 & 0 & -1 & 0
\end{array}\right) \quad \ldots
$$

and it is also easy to see that these are the only ones.
(6) Any matrix satisfying (4) can after a change of basis be written in the form given in (5).

Proof. Let $x$ have norm 1 and let $y=A x$. (a) and (b) imply $x \cdot y=0$ and $|y|=1$. Using (a) twice more gives

$$
0=(x+y) \cdot A(x+y)=y \cdot A x+x \cdot A y
$$

so

$$
x \cdot A y=-y \cdot y=-1
$$

and since $|A y|=|y|=1$ it follows that $A y=-x$.
The last result shows that the behavior observed by Stein and Weiss is very rare among martingale transforms and in fact distinguishes "conjugation" and its generalizations to $\mathrm{R}^{2 \mathrm{n}}$ (the Hilbert transforms of Varopoulos (1980)) from the other martingale transforms. Faced with this situation, if we want to prove something for more general matrices we have to settle for something less than the conclusion of (4). The next result shows that we can weaken the condition on the matrices quite a bit without sacrificing too much in the conclusion.
(7) If $A$ has no real eigenvalue then there are $C$ and $\gamma$ which depend on $A$ and $P(B)$ so that $P\left(\sup _{t}\left|\left(A * 1_{B}\right)_{t}\right|>y\right) \geq C e^{-\gamma y}$

Before proving this we would like to make two remarks which explain the condition and the conclusion.

1. The result is false if $A$ has a real eigenvalue for if $V R^{d}$ is an associated real eigenvector and we

$$
\begin{aligned}
\text { let } y_{t} & =\frac{1}{2}+\int_{0}^{t} v \cdot d B_{s} \\
\text { let } \sigma & =\inf \left\{t: Y_{t} \notin(0,1)\right\} \\
\text { and let } X_{t} & =Y_{t \wedge \sigma}
\end{aligned}
$$

then $X_{\infty}=1\left(Y_{\sigma}=1\right)$ but

$$
\begin{aligned}
(A * X) & =\int_{0}^{\sigma} A v \cdot d B_{s} \\
& =\lambda \int_{0}^{\sigma} v \cdot d B_{s}=\lambda\left(Y_{\sigma}-Y_{0}\right)
\end{aligned}
$$

so $A * X$ is bounded.
2. Well known formulas for Brownian motion show that when $A=\left(\begin{array}{rr}0 & -1 \\ 1 & 0\end{array}\right)$ the left hand side of (7) is $\sim \mathrm{Ce}^{-\gamma y}$ (here $C, \gamma \quad(0, \infty)$ are constants whose values may change from line to line) and the John Nirenberg inequality (see Meyer (1976) or Durrett (1984), Section 7.6) implies that for any matrix A

$$
P\left(\sup _{t}\left|\left(A * 1_{B}\right)_{t}\right|>y\right) \leq C e^{-\gamma y}
$$

where $C, \gamma$ depend only on $A$ and $P(B)$ so we cannot hope for a better lower bound.

Proof of (7). Let $X=1_{B}, \quad X_{t}=E\left(X \mid \exists_{t}\right)$. We will prove (7) by showing that although $\left(X_{t},(A * X)_{t}\right)$ may not be a time change of Brownian motion, there is a part of $A * X$ which is independent of $X$ and which is a time change of a Brownian mition run for an amount of time $\geq \varepsilon T$ where $T$ is the time defined in the proof of (2).

To isolate the part of $A * X$ we want, we introduce the following orthogonal decomposition of Ax

$$
A x=C(x) x+F(x)
$$

where $C(x)$ is a number and $F(x) \in R^{d}$ has

$$
F(x) \cdot x+0
$$

It is easy to see that the last two equations specify $C(x)$ and $F(x)$ and we have $|F(x)| \leq|A x|$. To prove (7) we need a bound in the other direction. To do this we observe that if $A$ has no real eigenvalues then $F(x) \neq 0$ for all $x \neq 0$ and scaling implies that for $y \neq 0$

$$
|F(y)|=|y|\left|F\left(\frac{y}{|y|}\right)\right|
$$

so we have
(8)

$$
\inf _{y \neq 0} \frac{|F(y)|}{|y|}=\inf _{z,|z|=1}|F(z)|>0
$$

With (8) established our next step is to decompose $(A * X)_{t}$. If

$$
x_{t}=\int_{0}^{t} H_{s} \cdot d B_{s}
$$

then

$$
(A * X)=\int\left(A H_{S}\right) \cdot d B_{S}
$$

so we let

$$
\begin{aligned}
& Y_{t}=\int_{0}^{t} c\left(H_{s}\right) H_{s} \cdot d B_{s} \\
& Z_{t}=\int_{0}^{t} F\left(H_{s}\right) \cdot d B_{s} .
\end{aligned}
$$

The formula for the covariance of two stochastic integrals (see Meyer (1976) or Durrett (1984) Chapter 2) implies

$$
\langle x, z\rangle_{t}=\int_{0}^{t} F\left(H_{s}\right) \cdot H_{s} d s=0
$$

and (8) tells us that

$$
\begin{aligned}
\langle z\rangle_{t} & \equiv\langle z, z\rangle_{t}=\int_{0}^{t}\left|F\left(H_{s}\right)\right|^{2} d s \\
& \geq \varepsilon^{2}\langle x\rangle_{t}
\end{aligned}
$$

At this point we have found the part of $A * X$ we referred to at the beginning of the proof. The next step is to show $Z$ has the desired properties. To do this we let

$$
\left.\gamma(u)=\inf \left\{t:\langle z\rangle_{t}\right\rangle u\right\} \text { for } u<\langle z\rangle_{\infty}
$$

and define

$$
W_{u}= \begin{cases}Z_{\gamma}(u) & u<\langle Z\rangle_{\infty} \\ Z_{\infty}+\hat{B}_{u-Z_{\infty}} & u \geq\langle Z\rangle_{\infty}\end{cases}
$$

where $\hat{B}$ is a one dimensional Brownian motion which is independent of the d-dimensional Brownian motion $B$. We have added $\hat{B}$ after the end of $Z$ so that the following holds.
(9) $W$ is a Brownian motion which is independent of $\sigma\left(X_{t}, t \geq 0\right)$.

Proof. This is a consequence of a theorem of F. Knight (1971) but the proof is short so we will prove it. It is easy to check that $W_{u} u \geq 0$ is a local martingale and $W_{u} \equiv u$ (for more details see Meyer (1976), or Durrett (1984)

Section 2.11) so it follows from Levy's characterization that $w_{u}$ is a Brownian motion. To check the independence
let

$$
u=\int f_{s} d x_{s}
$$

and let

$$
v=\int g_{s} d W_{s}
$$

be stochastic integrals with

$$
\left.\int\left|f_{s}\right|^{2} d s x\right\rangle_{s}, \int\left|g_{s}\right|^{2} d s<\infty
$$

and $g_{s}=0$ for $s \geq\langle Z\rangle_{\infty}$. Unscrambling the definitions we see that

$$
\int f_{s} d X_{s}=\int\left(f_{s} H_{s}\right) \cdot d B_{s}
$$

and

$$
\begin{aligned}
\int g_{s} d W_{s} & =\int g_{s} d Z_{\gamma(s)} \\
& =\int g\left(\langle Z\rangle_{t}\right) d Z_{t} \\
& =\int g\left(\langle Z\rangle_{s}\right) F\left(H_{s}\right) \cdot d B_{s},
\end{aligned}
$$

so it follows from the formula for the covariance of two stochastic integrals that

$$
E U V=E \int f_{s} g\left(\langle Z\rangle_{s}\right) H_{s} \cdot F\left(H_{s}\right) d s=0 .
$$

It is trivial that we have $E U V=0$ if $g_{s}=0$ for $s \leq\langle Z\rangle_{\infty}$ so the last equality holds for any $f, g$ which satisfy (*) and hence for any $V \in L^{2}\left(\sigma\left(X_{t}\right.\right.$ : $t \geq 0)$ ) and $V \in L^{2}\left(\sigma\left(W_{u}: u \geq 0\right)\right)$ which proves (9).

With (9) established the rest is simple but requires a little trickery. $Z_{t}$ is a time change of $W_{u} u\left\langle\langle Z\rangle_{\infty}\right.$ and by (8) $\langle Z\rangle_{\infty} \geq \varepsilon^{2}\langle X\rangle_{\infty}$, so we have

$$
\sup _{t}\left|z_{t}\right| \geq \sup _{u \leq \varepsilon}^{2}\langle x\rangle_{\infty}\left|w_{u}\right| .
$$

where $\langle x\rangle_{\infty} \in \sigma\left(X_{t}: t \geq 0\right)$ is independent of $W$. To get a lower bound on $\sup _{t}\left|(A * X)_{t}\right|$ find the first point $u_{0} \leq \varepsilon^{2}\langle X\rangle_{\infty}$ where the sup on the right occurs. If we let $t_{0}=\gamma\left(u_{0}\right)$ (which is finite since $u_{0} \leq \varepsilon^{2}\langle X\rangle_{\infty} \leq\langle Z\rangle_{\infty}$ then

$$
(A * X)_{t_{0}}=Y_{t_{0}}+Z_{\gamma\left(u_{0}\right)} .
$$

At this point we could get unlucky and $Y_{t_{0}}$ could cancel $Z_{\gamma}\left(u_{0}\right)$, but the sign of $Z_{\gamma\left(u_{0}\right)}$ is independent of the sign of $\gamma_{t_{0}}$ so at least $\frac{1}{2}$ of the time $Y_{t_{0}}$ will make $\left|(A * X)_{t_{0}}\right| \geq\left|Z_{\gamma\left(u_{0}\right)}\right|$ and it follows that

$$
P\left(\left|(A * X)_{t_{0}}\right|>y\right) \geq \frac{1}{2} P\left(\sup _{u \leq \varepsilon}^{2}\langle x\rangle_{\infty}\left|W_{u}\right|>y\right)
$$

To compute the quantity on the right hand side and complete the proof of (8) we observe that since $X_{\infty}$ is independent of $W_{u}$

and the distribution of the right hand side is given in Remark 2.
Having proved (7) for one matrix, it is natural, especially if you have heard of Janson's (1977) theorem (see e.g. Durrett (1984), Section 6.7.), to ask what happens if we have a family of matrices without a common real eigenvector. The answer is just what you should expect:
(10) If $A^{1}, \ldots, A^{m}$ are matrices without a common real eigenvector then there are constants $C$ and $\gamma$ which only depend on $P(B)$ (and the matrices) so that

$$
P\left(\sup _{i} \sup _{t}\left|\left(A^{i} * 1_{B}\right)\right| \quad y\right) \geq C e^{-\gamma} y
$$

Since this is a rather straightforward generalization of (7) we will explain why we want to prove this before we describe how to do it. In our discussion of (4) above we observed that if $d$ is odd then $A$ must have a real eigenvector, so the hypothesis of (7) cannot be satisfied in this case. With the matrices of (5) in mind you might realize that in the first nontrivial case ( $d=3$ ) it is easy to write down two matrices which have no common real eigenvector

$$
A_{1}=\left(\begin{array}{rrr}
0 & 1 & 0 \\
-1 & 0 & 0 \\
0 & 0 & 0
\end{array}\right) \quad A_{2}=\left(\begin{array}{rrr}
0 & 0 & 1 \\
0 & 0 & 0 \\
-1 & 0 & 0
\end{array}\right)
$$

Congratulations, you have just (re)discovered the Riesz transforms. Gundy and Varopolous (1979) (and later by a different method Gundy and Silverstein (1982)) have shown that if we define a process $W_{t}-\infty<t \leq 0$ in $H=R^{n} \times(0, \infty)$ which is a Brownian motion which "starts at time $-\infty$ from Lebesgue measure on $R^{n} \times\{\infty\}$ and exits $H$ at time $0, "$ then the Reisz transforms are related to martingale transforms of $W$.

To explain the relationship we need some notation. Let $f$ be a function on $\partial \mathrm{H}$ which is in $\mathrm{L}^{2}$ and let

$$
u(z)=E_{z} f\left(B_{\tau}\right)
$$

where $\tau=\inf \left\{t: B_{t} \notin H\right\}$. If we let $A^{i}$ be the matrix which has

$$
A_{j k}^{i}=\left\{\begin{array}{rl}
1 & j=1 \quad k=\mathbf{i} \\
-1 & j=\mathbf{i} \quad k=1 \\
0 & \text { otherwise }
\end{array}\right.
$$

then the ith Riesz transform may be written as

$$
R_{i} f\left(w_{0}\right)=E\left(\int_{-\infty}^{0} A^{i} \nabla u\left(w_{s}\right) \cdot d w_{s} \mid w_{0}\right)
$$

Since the Riesz transforms are (for the theory of Hardy spaces at least) the appropriate generalization to $H$ of conjugation in $D$, it is natural to ask if

$$
\left|\left\{x: \sup _{j} R_{i} 1_{B}(x, 0)>\lambda\right\}\right| \geq C e^{-\gamma \lambda}
$$

where $C$ and $\gamma$ are constants which depend only on $|B|$. (10) shows that the analogous result is true for martingale transforms and that the stochastic integral in (11) is unbounded. Unfortunately the conditional expectation might convert the integral into a bounded function so we have not been able to use this to solve the (still open) question posed above.

Proof. For simplicity, we will give the proof only for $m=2$. The reader can obtain a proof of the general result by changing 2 to $m$ and inserting ... at appropriate points. As in the proof of (7) we begin by introducing orthogonal decompositions

$$
\begin{aligned}
& A^{1} x=c^{1}(x) x+F^{1}(x) \\
& A^{2} x=c_{0}^{2}(x) x+c_{1}^{2}(x) F^{1}(x)+F^{2}(x)
\end{aligned}
$$

where $F^{i}(x) \cdot x=0 \quad i=1,2$ and $F^{1}(x) \cdot F^{2}(x)=0$.
Now if $A^{1}$ and $A^{2}$ have no common real eigenvector then

$$
\left\{F^{1}(x)=0\right\} \cap\left\{c_{1}^{2}(x) F^{1}(x)+F^{2}(x)=0\right\}=\emptyset
$$

i.e. $\left\{F^{1}(x)=0\right\} \cap\left\{F^{2}(x)=0\right\}=\emptyset$ and repeating the proof of (8) shows

$$
\begin{equation*}
\inf _{x \neq 0} \frac{\left|F^{1}(x)+F^{2}(x)\right|}{|x|} \equiv \varepsilon>0 \tag{11}
\end{equation*}
$$

The next step is the decompose the $\left(A^{i} * X\right)_{t}$ and time change some of the pieces to produce independent Brownian motions.

Let $Z_{t}^{i}=\int_{0}^{t} F^{i}\left(B_{s}\right) \cdot d B_{s}$
let $Y_{t}^{i}=\left(A^{i} * X\right)_{t}-Z_{t}^{i}$
let $\left.\gamma_{i}(u)=\inf \left\{t:\langle Z\rangle_{t}\right\rangle u\right\}$
and let

$$
W_{u}^{i}= \begin{cases}z_{\gamma_{i}}^{i}(u) & u\left\langle\left\langle z^{i}\right\rangle_{\infty}\right. \\ z_{\infty}^{i}+\hat{B}_{u-\left\langle z^{i}\right\rangle_{\infty}}^{i} & u \geq\left\langle z^{i}\right\rangle_{\infty}\end{cases}
$$

where $\hat{B}^{1}$ and $\hat{B}^{2}$ are independent Brownian motions which are independent of $B$.
A simple generalization of (9) (or invoking Knight's theorem) implies that $W_{u}^{1}$ and $W_{u}^{2}$ are independent Brownian motions which are independent of $\sigma\left(x_{t}: t \geq 0\right)$ and (11) implies that $\left\langle Z^{1}\right\rangle_{\infty}+\left\langle z^{2}\right\rangle_{\infty} \geq \varepsilon^{2}\langle x\rangle_{\infty}$ so now we can complete the proof almost as before.

Let $j=, \inf \left\{i:\left\langle z^{i}\right\rangle_{\infty} \geq \varepsilon^{2} / 2\langle x\rangle_{\infty}\right\}$
Let $u_{0}$ be the first point at which

$$
\sup _{.} u \leq\left(\varepsilon^{2} / 2\right)\langle x\rangle_{\infty}\left|W_{u}^{j}\right| \text { is attained. }
$$

Let $t_{0}=\gamma^{j}\left(u_{0}\right)$

$$
\left(A^{j} * X\right)_{t_{0}}=Y_{t_{0}}^{j}+Z_{\gamma^{j}}^{j}\left(u_{0}\right)
$$

and again the signs of the two terms on the right are independent so

$$
P\left(\left|\left(A^{j} * x\right)_{t_{0}}\right|>y\right) \geq \frac{1}{2} P\left(\sup _{u \leq\left(\varepsilon^{2} / 2\right)\langle x\rangle_{\infty}}\left|W_{u}^{j}\right|>y\right)
$$

proving the desired result.

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